

Symmetries of \mathbb{R}^n

$$\text{Isom}(\mathbb{R}^n)$$

$$\mathbb{R}^n < \text{Isom}(\mathbb{R}^n)$$

$$v \longleftrightarrow t_v$$

$$\text{Isom}(\mathbb{R}^n) \xrightarrow{\alpha} O_n(\mathbb{R})$$

surjective

1st isom thm

↓

$$\text{Isom}(\mathbb{R}^n) / \ker \alpha \cong O_n(\mathbb{R})$$

$$\frac{\text{Isom}(\mathbb{R}^n)}{\mathbb{R}^n} \cong O_n(\mathbb{R})$$

$$O_n(\mathbb{R}) \xrightarrow{\det} \pm 1$$

$\ker(\det)$ "orientation preserving"

" $SO_n(\mathbb{R})$ "

$$\det = -1$$

const. of $SO_n(\mathbb{R})$

$$n=2$$

$$SO_2(\mathbb{R})$$

rotations

$$\begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}$$

orientation rev: reflectors

$n=3$

$SO_3(\mathbb{R})$

rotations

reflection free

$T \in M_n(\mathbb{R})$ n odd then T has an eigenvector.

$$\chi_T(x) = \det(xI_n - T) \text{ degree } n.$$

$x^n + \dots$

has a root λ

$$(\lambda I - T)v = 0$$

singular \Downarrow

v an eigen vector.

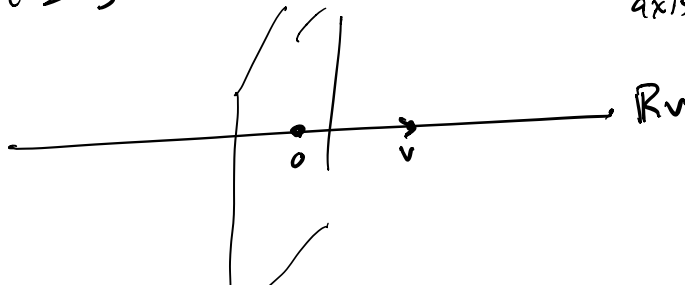
if $T \in O_3(\mathbb{R})$

$$Tv = \lambda v \Rightarrow \lambda = \pm 1$$

$$\|v\| = \|Tv\| = |\lambda| \|v\|$$

$$= |\lambda| \|v\| \Rightarrow |\lambda| = 1$$

if $T \in SO_3(\mathbb{R})$ then T is a rotation through some axis



So T preserves plane $P \subset \mathbb{R}^3$ perp to v .

i.e. T gives an isometry $T|_P$ of P

$\Rightarrow T|_P$ either rotation or reflection

T as a matrix has block-diagonal form

$$T = \left[\begin{array}{c|c} T|_P & 0 \\ \hline 0 & \pm 1 \end{array} \right] \quad \begin{array}{l} \text{basis } \underbrace{b_1, b_2}_{\substack{\text{basis for} \\ P}} \\ \uparrow \\ v \\ \text{e. vect.} \end{array}$$

if $\det = 1$

either $T|_P$ an ± 1 has $\det = 1$

or $T|_P$ is ± 1 has $\det = -1$

$\Rightarrow T|_P$ is a rotation, $\ell = \mathbb{R}v$ is fixed

or $T|_P$ is a reflection, -1 in corner

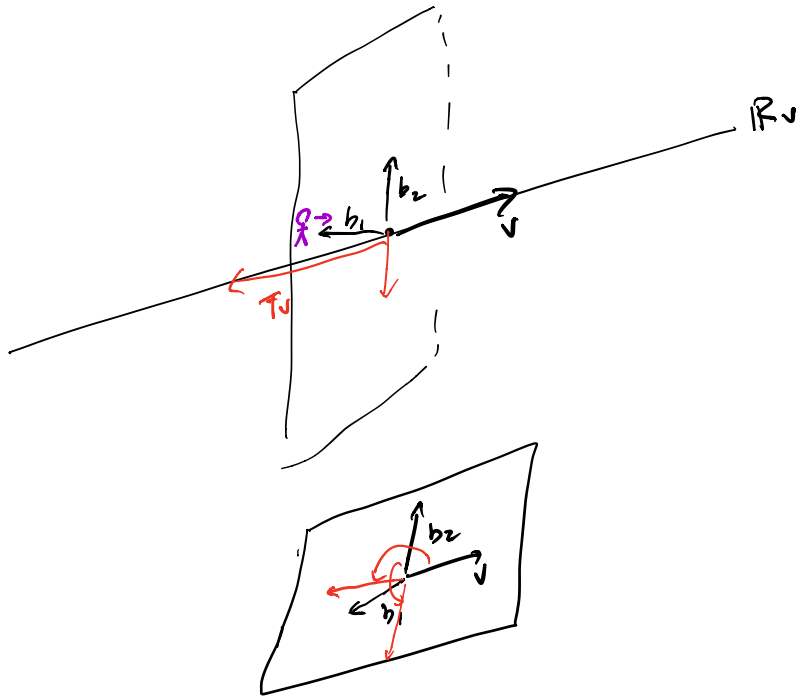
can find b_1, b_2 an orth basis for P w/

$$Tb_1 = b_1$$

$$Tb_2 = -b_2$$

$$\Rightarrow T = \begin{bmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{bmatrix}$$

180° degree rotation
in $\langle b_2, v \rangle$
plane.



Def $H = \{a + bi + cj + dk \mid a, b, c, d \in \mathbb{R}\}$
 $i^2 = j^2 = k^2 = -1$
 $ij = k = -ji$

Quaternions

$$q = a + bi + cj + dk \quad \bar{q} = a - bi - cj - dk$$

$$\|q\|^2 = q\bar{q} = a^2 + b^2 + c^2 + d^2 \in \mathbb{R}$$

$$H^\times = \{q \in H \mid q \neq 0\}$$

group under mult.

$$q \cdot \frac{\bar{q}}{\|q\|^2} = 1$$

$$q^{-1} = \frac{\bar{q}}{\|q\|^2}$$

$\mathbb{R}^3 \longleftrightarrow$ "pure quaternions"

$$= \{bi+cj+dk \mid b, c, d \in \mathbb{R}\} \subset \mathbb{H}$$

and then if $v \in \mathbb{R}^3 \subset \mathbb{H}$ and $g \in \mathbb{H}^*$

map $v \mapsto gv g^{-1}$ describes a rotation in \mathbb{R}^3

get a map $\mathbb{H}^x \xrightarrow{\varphi} SO_3(\mathbb{R})$

$$g \mapsto (v \mapsto gv g^{-1})$$

if $g \in \mathbb{R}^x \subset \mathbb{H}^x$ then $g \in \ker \varphi$

$$\text{Spin}_3(\mathbb{R}) = \mathbb{H}^x / \mathbb{R}^x \longrightarrow SO_3(\mathbb{R})$$

$$\tilde{\mathbb{R}}_1 \longrightarrow \mathbb{R}$$

$$\tilde{\mathbb{R}}_2 \longrightarrow$$

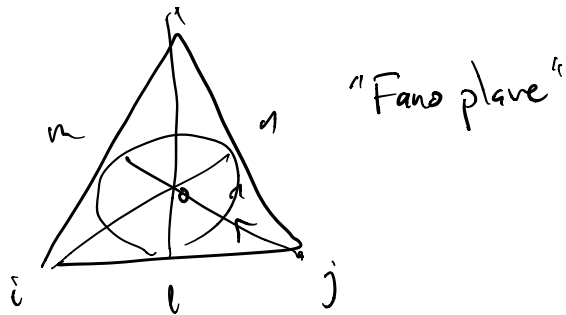
$$gv g^{-1} = v g g^{-1} = v$$

if $g = \lambda \in \mathbb{R}^x \subset \mathbb{H}^x$

$$g \cdot p = p \cdot g$$

$$\mathbb{D} = \{a + bi + cj + dk + el + fm + gn + ho\}$$

$$i^2 = j^2 = \dots = -1$$



Theorem (Frobenius 1880's)

if A is a division algebra of finite dimension over \mathbb{R}

then $A = \mathbb{R}, \mathbb{C}, \mathbb{H}, \mathbb{O}$.

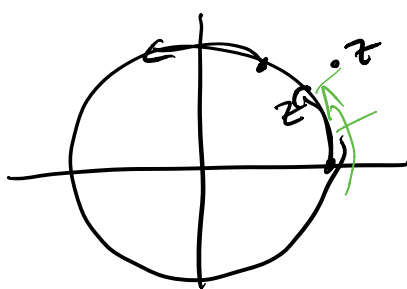
↑
not associative

Alternate approach to rotations in higher dim

Clifford algebras

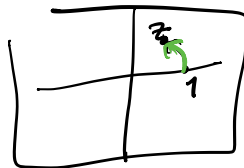
only reasonable ways to describe Spin phenomena
in higher dim.

$$wz = w \Rightarrow z = 1$$



choose $z \neq 1$ in \mathbb{C}

$$\|z\| = 1$$





Thm 1960 Adams
Tangent bundles for spheres are only parallelizable
in dims 1, 3, 7.